Shifts on a Deformed Hilbert Space

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In this paper we observe that the creation operator on a deformed Hilbert space is the product of an ordinary shift and a diagonal operator.

1. INTRODUCTION

We consider a Hilbert space which is spanned by the vectors $\{e_n, n =$ 0, 1, 2, . . . }. These vectors are generated by the action of the *creation operator* a^* on the vector e_0 which we call a *vacuum* vector in the Hilbert space. The hermitian conjugate of *a** is the *annihilation operator a*. Together they satisfy the following relations:

$$
aa^* = qa^*a + 1
$$

\n
$$
\langle e_0, e_0 \rangle = 1
$$

\n
$$
e_n = (a^*)^n e_0
$$

\n
$$
ae_0 = 0
$$

\n(1)

We call this Hilbert space a *deformed Hilbert space* and denote it by H_q , where *q* is a deforming parameter ranging over $0 \le q \le 1$. This space was discussed in ref. 1.

Using the relations in (1), we observe the following:

$$
a^*e_n = e_{n+1}
$$

\n
$$
ae_n = [n]e_{n-1}
$$
\n(2)

where we take $[n] = (1 - q^n)/(1 - q) = 1 + q + q^2 + \ldots + q^{n-1}$.

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47

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To prove $\langle e_m, e_n \rangle = [n]! \delta_{nm}$ we proceed to prove $\langle e_n, e_n \rangle = [n]!$:

$$
||e_{n+1}||^2 = \langle a^* e_n, e_{n+1} \rangle
$$

= $\langle e_n, a e_{n+1} \rangle$
= $\langle e_n, [n+1]e_n \rangle$
= $[n+1]||e_n||^2$

Hence, $||e_n||^2 = [n] ||e_{n-1}||^2 = [n] [n-1] \dots [2] [1] ||e_0||^2 = [n]! ||e_0||^2$. Since by (1), $||e_0||^2 = 1$, we have

$$
\langle e_n, e_n \rangle = ||e_n||^2 = [n]!
$$

Now,

$$
\langle e_n, e_m \rangle = \langle a^{*n} e_0, a^{*m} e_0 \rangle = \langle e_0, a^n a^{*m} e_0 \rangle
$$

= 0 if $m > n$ or $m < n$

Thus,

$$
\langle e_n, e_m \rangle = [n]!\,\delta_{nm}
$$

Hence, the following can be derived using relations (1):

$$
a^* e_n = e_{n+1}
$$

\n
$$
ae_n = [n]e_{n-1}
$$

\n
$$
\langle e_n, e_m \rangle = [n]! \delta_{nm}
$$
\n(3)

where $[n]! = [n] \dots [2][1]$; $[0]! = 1$.

The vectors $\{([n]!)^{-1/2} e_n\}$ form an orthonormal basis and the Hilbert space H_q consists of all vectors $f = \sum_{n=0}^{\infty} f_n e_n$ with complex f_n such that

$$
\langle f, f \rangle = \sum_{n=0}^{\infty} |f_n|^2[n]!
$$

is finite.

If $g = \sum_{n=0}^{\infty} g_n e_n$ is also a vector in the space, then

$$
\langle f, g \rangle = \sum_{n=0}^{\infty} \bar{f}_n g_n[n]!
$$

where the bar denotes the complex conjugate.

In this paper we prove two theorems. The content of the first theorem is that the shift operator is always a bounded operator in H_q . The content of the second theorem is that the shift operator on H_q is unitarily equivalent to

the weighted shift on l^2 . Using these theorems, we have shown that the creation operator mentioned above is the product of an ordinary shift and a diagonal operator [2].

2. THEOREMS

Theorem 1. The shift is always an operator in H_q , that is, if $f = \sum_{n=0}^{\infty}$ $f_n e_n = \{f_0, f_1, f_2, \ldots\} \in H_q$, then $Sf = \{0, f_0, f_1, \ldots\} \in H_q$ and as *f* varies over H_q , ||Sf|| is bounded by a constant multiple of ||f||.

Proof. It is necessary that $||e_{n+1}|| \le \alpha ||e_n||$, where e_n is the vector whose coordinate with index *n* is 1 and all other coordinates are 0. Since

$$
||e_n||^2 = \langle e_n, e_n \rangle = [n]! \equiv p_n
$$

this condition says that

$$
\frac{p_{n+1}}{p_n} = \frac{[n+1]!}{[n]!} = [n+1] = \frac{1-q^{n+1}}{1-q} \to \frac{1}{1-q} \equiv \alpha
$$

as $n \to \infty$.

Thus the condition $||e_{n+1}|| \le \alpha ||e_n||$ says that the sequence $\{p_{n+1}/p_n\}$ is bounded with bound $1/(1 - q)$.

This condition is also sufficient. For, if $p_{n+1}/p_n \leq 1/(1 - q)$ for all *n*, then

$$
||Sf||^2 = \sum_{n=1}^{\infty} p_n |f_{n-1}|^2
$$

=
$$
\sum_{n=1}^{\infty} \frac{p_n}{p_{n-1}} p_{n-1} |f_{n-1}|^2
$$

$$
\leq \frac{1}{1-q} \sum_{n=0}^{\infty} p_n |f_n|^2
$$

=
$$
\frac{1}{1-q} ||f||^2
$$

Hence

$$
||Sf|| \le \frac{1}{(1-q)^{1/2}} ||f||
$$

Theorem 2. The shift S on H_q is unitarily equivalent to the weighted shift *T*, with weights $\{\sqrt{[n+1]}\}\$, on *l*².

Proof. We observe that $\{[n]!\}$ is a sequence of positive numbers with $\{[n + 1]!/[n]!\}$ a bounded sequence.

If
$$
f = \{f_0, f_1, f_2, ...\} \in H_q
$$
, we write
\n $Uf = \{1.f_0, \sqrt{[1]!}f_1, \sqrt{[2]!}f_2, ...\}$. Then,
\n
$$
||Uf||^2 = ||f||^2 = \sum_{n=0}^{\infty} [n]![f_n]^2 < \infty
$$

Thus, $Uf \in l^2$. Hence, U maps H_q into l^2 linearly.

If $g = \{g_0, g_1, g_2, ...\} \in l^2$ and if $f_n = g_n/\sqrt{[n]!}$, then $\Sigma_{n=0}^{\infty} [n]! |f_n|^2$
= $\Sigma_{n=0}^{\infty} |g_n|^2 < \infty$. This proves that U maps H_q onto l^2 . Now,

$$
USU^{-1}{g_0, g_1, g_2, \ldots} = US\left\{\frac{g_0}{\sqrt{[0]!}}, \frac{g_1}{\sqrt{[1]!}}, \frac{g_2}{\sqrt{[2]!}}, \ldots\right\}
$$

= $U\left\{0, \frac{g_0}{\sqrt{[0]!}}, \frac{g_1}{\sqrt{[1]!}}, \frac{g_2}{\sqrt{[2]!}}, \ldots\right\}$
= $\left\{0, \sqrt{\frac{[1]!}{[0]!}} g_0, \sqrt{\frac{[2]!}{[1]!}} g_1, \sqrt{\frac{[3]!}{[2]!}} g_2, \ldots\right\}$
= $T{g_0, g_1, g_2, \ldots}$

Thus, U transforms S onto T . That is, the transform of the ordinary shift on H_q is a weighted shift on l^2 .

Observation. Because of Theorem 2, a shift S on H_q is the weighted shift T, with weights $\{\sqrt{[n+1]}\}\$ on l^2 ; we have $T = U\dot{D}$, where U is an ordinary shift and D is a diagonal operator with diagonals $\{[n + 1]!/[n]!\}$. We see that

$$
a^* = UD
$$

and

$$
ae_n = D^*U^*e_n = D^*e_{n-1} = [n]e_{n-1}
$$

We also see that

$$
\|a\| = \|a^*\| = \sup_n \sqrt{[n+1]} = \sup_n \sqrt{\frac{1 - q^{n+1}}{1 - q}} = (1 - q)^{-1/2}
$$

REFERENCES

- 1. Arik, M., and Coon, D. D., Hilbert spaces of analytic functions and generalized coherent states, JMP 17, 524-527 (1976).
- 2. Halmos, P. R., A Hilbert Space Problem Book, Springer-Verlag, Berlin (1982).